

Lecture # 23

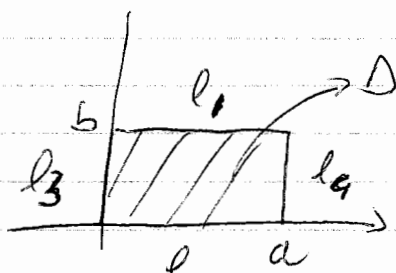
Special Domains and separation of variables

For certain domains one could solve the Laplace equation using separation of variables

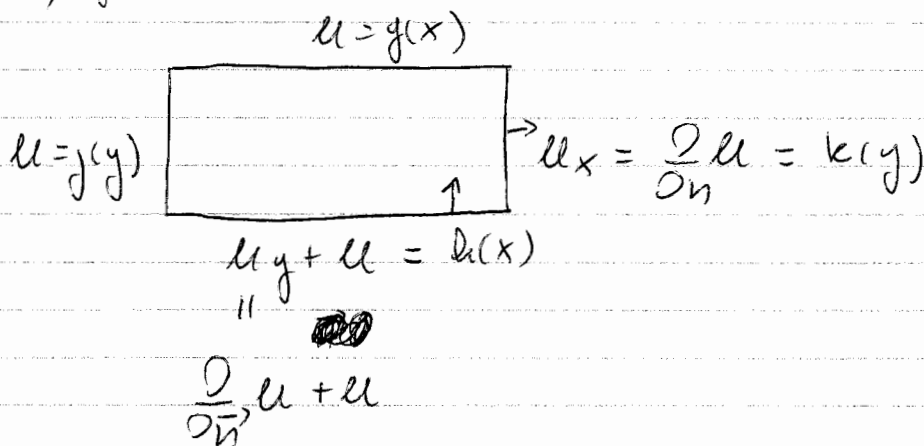
Ex 1: 2-D rectangle

$$\Delta_2 u = u_{xx} + u_{yy} = 0$$

$$D = [0, a) \times (0, b)$$



On each of the sides one can put either Dirichlet, or Neuman or Robin conditions



$u$  solves  $\begin{cases} \Delta_2 u = 0 \\ u|_{\partial D} = (g, h, j, k) \end{cases} \implies u = u_1 + u_2 + u_3 + u_4$

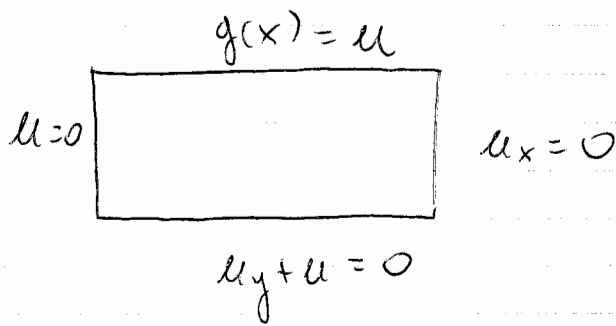
$u_i$

$$\begin{cases} u_1 = 0 \\ u_1|_{\partial D} = (g, 0, 0, 0) \end{cases} \begin{cases} \Delta u_2 = 0 \\ u_2|_{\partial D} = (0, h, 0, 0) \end{cases} \begin{cases} \Delta u_3 = 0 \\ u_3|_{\partial D} = (0, 0, j, 0) \end{cases} \begin{cases} \Delta u_4 = 0 \\ u_4|_{\partial D} = (0, 0, 0, k) \end{cases}$$

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Then the ~~to~~ problems below can be solved by separation of variables in the way discussed before.

In the book is discussed the case



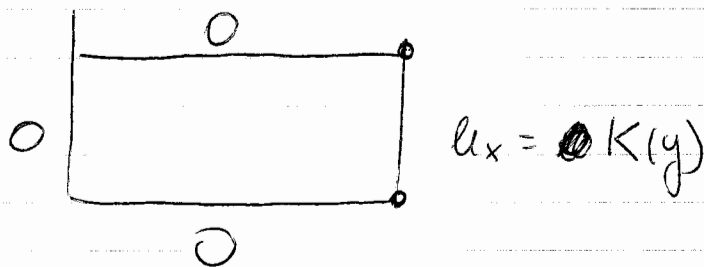
~~the~~  $u(x, y) = \sum_{n=0}^{\infty} A_n \sin \beta_n x (\beta_n \cosh \beta_n y - \sinh \beta_n y)$

and  $\beta_n = (n + \frac{1}{2}) \frac{\pi}{a} \quad n = 0, 1, \dots$

and  $A_n$  s.t.  $g(x) = \sum_{n=0}^{\infty} A_n (\beta_n \cosh \beta_n b - \sinh \beta_n b) \sin \beta_n x$

$0 < x < a$

Here let's solve  $u_4$



Use separate variables  $u(x, y) = X(x) Y(y)$

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$$\Delta_2 X(x) Y(y) = X''(x) Y(y) + X(x) Y''(y) = 0$$

$$\frac{X''}{X} + \frac{Y''}{Y} = 0$$

$$\begin{cases} \frac{X''}{X} = +\lambda & \frac{Y''}{Y} = -\lambda \\ 0 \leq x \leq a & 0 \leq y \leq b \end{cases}$$

In particular from the data in  $I_1$  and  $I_2$  it follows

$$\begin{cases} Y'' = \lambda Y \\ Y'(0) + Y(0) = 0 \\ Y(b) = 0 \end{cases}$$

this is an eigenvalue problem

~~with~~ ~~Robin~~ ~~conditions~~ ~~at~~ ~~the~~ ~~boundary~~ ~~points~~ ~~0~~ ~~and~~ ~~b~~

with Robin' conditions at

~~the~~ ~~boundary~~ ~~points~~

0 and b

Assume  $\lambda_0, \lambda_1, \dots, \lambda_n, \dots$  are the eigenvalues then  
 using the notation in the book for RC:

$$Y'(0) - a_0 Y = 0 \Rightarrow a_0 = -1$$

$$Y'(b) + a_1 Y = 0$$

Find eigenvalues  $\lambda_n$  and associated eigenfunctions  $Y_n$   
then write

$$X(y) = \sum A_n Y_n(y)$$

Once  $\lambda_n$  are obtained go back to equation of  $X$   
and solve

$$\begin{cases} X''_n = -\lambda_n X_n & \text{in } (0, a) \\ X_n(0) = 0 \end{cases}$$

then put everything together

$$u_2(x, y) = \sum_{n=0}^{\infty} A_n Y_n(y) (B_n X_n(x))$$

Finally use the fact that  $u_2(a, y) = h(x)$

to go back and determine  $A_n$  and  $B_n$ .

• A similar and lengthy construction solves also  
 $\Delta_3 u$  in  $D$  when  $D$  is a box.

### Dirichlet Problem in a disk

$$\begin{cases} \Delta_2 u = 0 & x^2 + y^2 \leq a^2 \\ u = h(\theta) & \text{in } x^2 + y^2 = a^2 \end{cases}$$

We already proved that

$$\Delta_2 u = u_{xx} + u_{yy} = u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta}$$

We use separation of variables  $u(r, \theta) = R(r) \Theta(\theta)$

$$\Theta R'' + \frac{1}{r} R \Theta' + \frac{1}{r^2} \Theta'' R = 0$$

$$\Theta (R'' + \frac{1}{r} R') + \Theta'' \frac{R}{r^2} = 0$$

divide by multiply by  $\frac{r^2}{R\Theta} \Rightarrow$

$$\frac{r^2 R'' + r R'}{R} + \frac{\Theta''}{\Theta} = 0$$

hence  $r^2 R'' + r R' - \lambda R = 0$

$$\Theta'' = -\lambda \Theta$$

clearly  $\Theta(\theta)$  is a  $2\pi$ -periodic function so we look at

$$\begin{cases} \Theta'' = -\lambda \Theta \\ \Theta(\theta) = \Theta(2\pi + \theta) \end{cases}$$

A simple calculation gives  $\lambda = n^2 \quad n = 1, 2, \dots$

$$\Theta(\theta) = A \cos n\theta + B \sin n\theta$$

$$\lambda = 0 \quad \Theta(\theta) = A$$

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Go back to  $R$ 

$$r^2 R'' + r R' - n^2 R = 0$$

we look for

$$R(r) = r^\alpha$$

$$r^2 \alpha(\alpha-1) r^{\alpha-2} + r \alpha r^{\alpha-1} - n^2 r^\alpha = 0$$

$$r^\alpha (\alpha(\alpha-1) + \alpha - n^2) = 0$$

$$\alpha^2 - \alpha + \alpha - n^2 = 0 \quad \alpha^2 = n^2 \quad \alpha = \pm n$$

$$R(r) = A r^n + B r^{-n}$$

$$u(r, \theta) = (C r^n + D r^{-n}) (A \cos n\theta + B \sin n\theta)$$

$$n = 1, 2, \dots$$

$$\text{For } n=0 \text{ we solve } r^2 R'' + r R' = 0$$

one can check that constants solve it, but also

$$\ln r \text{ since } (\ln r)' = \frac{1}{r} \quad (\ln r)'' = -\frac{1}{r^2}$$

$$\text{so we also have } u(r, \theta) = C + D \log r$$

Because we do not want singular solutions at  $0=r$   
 Since in the dish we have to have 2 derivatives, we

(2)

We discard the solutions  $r^n$  and  $r^{-n}$  so

$$\textcircled{A} u(r, \theta) = \frac{1}{2} A_0 + \sum_{n=1}^{\infty} r^n (A_n \cos n\theta + B_n \sin n\theta)$$

↓  
constant solution

by the boundary condition  $u(a, \theta) = h(\theta)$  we get

$$h(\theta) = \frac{1}{2} A_0 + \sum_{n=1}^{\infty} a^n (A_n \cos n\theta + B_n \sin n\theta)$$

$$\text{so } A_0 = \frac{1}{\pi} \int_0^{2\pi} h(s) ds$$

$$A_n = \frac{1}{a^n \pi} \int_0^{2\pi} h(s) \cos ns ds$$

$$B_n = \frac{1}{a^n \pi} \int_0^{2\pi} h(s) \sin ns ds$$

If we substitute this in  $\textcircled{A}$  we obtain

$$\begin{aligned} u(r, \theta) &= \int_0^{2\pi} h(s) \frac{ds}{2\pi} + \sum_{n=1}^{\infty} \frac{r^n}{a^n \pi} \int_0^{2\pi} h(s) \\ &\times \int_0^{2\pi} h(s) \left( \underbrace{\cos ns \cos n\theta + \sin ns \sin n\theta}_{\cos n(\theta-s)} \right) ds \\ &= \int_0^{2\pi} h(s) \left\{ 1 + 2 \sum_{n=1}^{\infty} \left(\frac{r}{a}\right)^n \cos n(\theta-s) \right\} \frac{ds}{2\pi} \end{aligned}$$

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Now in  $\{.3\}$  we find

$$1 + \sum_{n=1}^{\infty} \left(\frac{r}{a}\right)^n e^{in(\theta-s)} + \sum_{n=1}^{\infty} \left(\frac{r}{a}\right)^n e^{-in(\theta-s)}$$

$$= 1 + \sum_{n=1}^{\infty} z^n \quad \text{where } z = \frac{r}{a} e^{i(\theta-s)}$$

$$+ \sum_{n=1}^{\infty} \overline{z}^n \quad \text{and } |z| = \frac{r}{a} < 1 \quad \text{if } r < a$$

$$\text{so } = 1 + \frac{z}{1-z} + \frac{\overline{z}}{1-\overline{z}} =$$

$$= 1 + \frac{z - z\overline{z} + \overline{z} - z\overline{z}}{(1-z)(1-\overline{z})} = 1 + \frac{z + \overline{z} - 2z\overline{z}}{1 - \overline{z} - z + z\overline{z}}$$

$$= 1 + \frac{2\operatorname{Re}z - 2|z|^2}{1 - 2\operatorname{Re}z + |z|^2} = 1 + \frac{2\frac{r}{a} \cos(\theta-s) - 2\frac{r^2}{a^2}}{1 - 2\frac{r}{a} \cos(\theta-s) + \frac{r^2}{a^2}}$$

$$= 1 + \frac{2ar \cos(\theta-s) - r^2}{a^2 - r^2} = \frac{a^2 - 2ra \cos(\theta-s) + r^2}{a^2 - 2ra \cos(\theta-s) + r^2}$$

Hence

$$u(r, \theta) = (a^2 - r^2) \int_0^{2\pi} \frac{h(s)}{a^2 - 2ra \cos(\theta-s) + r^2} \frac{ds}{2\pi}$$

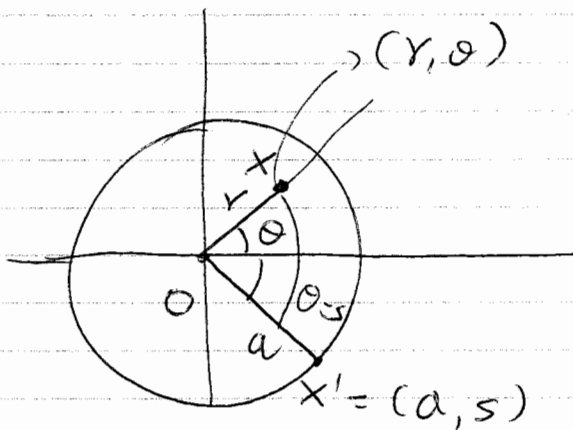
Poisson Formula



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Remark: The main point of this fact is that any harmonic function ~~can be written~~ on a disk can be written in terms of its value on the boundary of the disk!

Now let's go from polar coordinates back to ~~standard~~ Cartesian coordinates



$$\begin{aligned} \text{Notice that } |x - x'|^2 &= |(r \cos \theta, r \sin \theta) - (a \cos s, a \sin s)|^2 \\ &= |(r \cos \theta - a \cos s, r \sin \theta - a \sin s)|^2 \\ &= a^2 + r^2 - 2ar \cos(\theta - s) \end{aligned}$$

$$u(x) = \frac{a^2 - |x|^2}{a} \int_{|x'|=a} \frac{u(x')}{|x - x'|^2} d\sigma'$$

where  $\sigma' = a d\sigma = \text{arc length}$   $d\sigma' = a d\sigma$

## Consequences of Poisson Formula

(2) Maximum principle:

$u(x)$  harmonic function in  $D$  <sup>open</sup> connected domain

b) ~~then~~ there exists no  $x_0$  in  $D$  s.t.  $u(x_0) = \max_{\overline{D}} u$

or  $u(x_0) = \min_{\overline{D}} u$ , unless  $u = \text{const}$ .

Proof: ~~By~~ By contradiction ~~then~~ there is  $x_0$  <sup>in  $D$</sup>  s.t.

$$u(x_0) = \max_{\overline{D}} u = M \quad \text{so}$$

$$u(x) \leq u(x_0) = M \quad \text{for all } x \text{ in } \overline{D}$$

Because  $D$  is open there exists a disk centered at  $x_0$  and radius  $\varepsilon$ ,  $D_\varepsilon = D(x_0, \varepsilon)$ , all contained in  $D$ .

By the Poisson formula

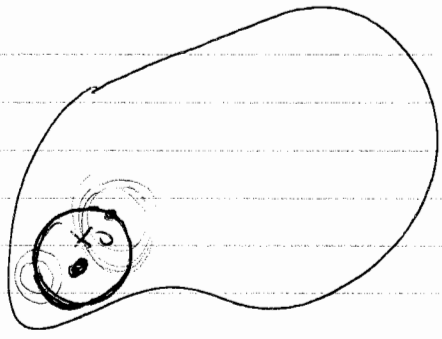
$$u(x_0) = \frac{\varepsilon^2 - |x_0|^2}{2\pi\varepsilon} \int_{|x'|=\varepsilon} \frac{u(x')}{|x_0 - x'|^2} d\sigma' \leq M$$

= average on  $D_\varepsilon$  ↓  
easy

so

$$M = u(x_0) = \text{average on } D_\varepsilon \leq M$$

If the average is the max then all points are max points so ~~on  $D_\varepsilon$~~   $u|_{\partial D_\varepsilon} = M$



By choosing the circles on the plane  
 you pick any point  $x$  inside  $D_\epsilon$ .  
 Then by taking a smaller radius  
 we can find a new disk ~~centered at~~  
~~at~~  $D(x_0, \tilde{\epsilon})$ ,  $\tilde{\epsilon} < \epsilon$  s.t.  $x$  is on the boundary  
 of  $D(x, \tilde{\epsilon})$ . By the above argument  $u(x) = u(x_0) = M$   
 so  $u = M$  on  $D_\epsilon$ . By covering  $D$  with such disks we  
 eventually find  $u \equiv M$  in  $D$ .

④ Mean Value problem

Assume  $x = 0$  then the Poisson formula reads

$$u(0) = \frac{a^2}{2\pi a} \int_{|x'|=a} \frac{u(x')}{a^2} d\sigma' = \text{Average of } u \text{ on circle of radius } a.$$

③ Differentiability: If  $u$  is a harmonic function  
 in any open set  $D$  of the plane, then  $u$  has all  
 the partial derivatives in any order.

Proof: Use Poisson formula:

Let  $x$  be in  $D$  on  $D_\epsilon = D(x, \epsilon)$  included in  $D$

~~max~~  $a^2$  then

$$u(x) = \frac{\epsilon^2 - |x|^2}{2\pi g} \int_{|x'| = \epsilon} \frac{u(x')}{|x - x'|^2} d\sigma'$$

Given the function  $u$  on integrating

$f(x, x') = \frac{u(x')}{|x - x'|^2}$  is continuous in  $x$  singular on  $|x'| = \epsilon$  and continuous in both variables

Then one can pass by derivative in side since  $\frac{1}{|x - x'|^2}$  holds all of them.