

Now consider the solution $Q(x,t)$ of the IVP

Lecture # 8

$$\begin{cases} u_t = k u_{xx} \\ u|_{t=0} = Q(x) = \begin{cases} 1 & \text{for } x > 0 \\ 0 & \text{for } x \leq 0 \end{cases} \end{cases}$$

If we dilate $Q(x) \rightarrow Q(\lambda x) = \begin{cases} 1 & \text{for } x > 0 \\ 0 & \text{for } x \leq 0 \end{cases} = Q(x)$

So we also expect it's solution not to change w.r.t. dilation, i.e. so that

$$Q(x,t) = Q(\lambda x, \lambda^2 t)$$

So we look for $Q(x,t) = g\left(\frac{x}{\sqrt{4kt}}\right)$
↓
normalization factor

Now by chain rule $\frac{x}{\sqrt{4kt}} = p \quad t > 0$

$$Q_t = \frac{dg}{dp} \cdot \frac{dp}{dt} = g'(p) \left(-\frac{1}{2t} \frac{x}{\sqrt{4kt}}\right)$$

$$Q_x = \frac{dg}{dp} \frac{dp}{dx} = \frac{g'(p)}{\sqrt{4kt}}$$

$$Q_{xx} = \frac{dQ_x}{dp} \frac{dp}{dx} = \frac{1}{4kt} g''(p)$$

$$0 = Q_t - Q_{xx} = \frac{1}{t} \left[-\frac{1}{2} p g'(p) - \frac{1}{4} g''(p) \right]$$

Then $g''(p) + 2p g'(p) = 0$

Now this is an ODE which solution is

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$$g(\beta) = c_1 \int_0^\beta e^{-\beta^2} d\beta + c_2$$

$$\text{so } Q(x,t) = c_1 \int_0^{\frac{x}{\sqrt{4kt}}} e^{-\beta^2} d\beta + c_2 \quad \forall t > 0$$

We need to check $\lim_{t \rightarrow 0}$

$$\text{if } x > 0 \quad \lim_{t \rightarrow 0^+} c_1 \int_0^{\frac{x}{\sqrt{4kt}}} e^{-\beta^2} d\beta + c_2$$

$$= c_1 \int_0^\infty e^{-\beta^2} d\beta + c_2$$

$$\text{if } x < 0 \quad \lim_{t \rightarrow 0^+} [\quad] = -c_1 \int_{-\infty}^0 e^{-\beta^2} d\beta + c_2$$

$$\text{but } \int_0^\infty e^{-\beta^2} d\beta = \frac{\sqrt{\pi}}{2} \text{ and } \int_{-\infty}^0 e^{-\beta^2} d\beta = -\frac{\sqrt{\pi}}{2}$$

$$\lim_{t \rightarrow 0^+} [\quad] = \begin{cases} c_1 \frac{\sqrt{\pi}}{2} + c_2 \times 1 = Q(x,0) = 1 & x > 0 \\ -c_1 \frac{\sqrt{\pi}}{2} + c_2 \times 0 = 0 & x < 0 \end{cases}$$

$$\begin{cases} c_1 \frac{\sqrt{\pi}}{2} + c_2 = 1 \\ -c_1 \frac{\sqrt{\pi}}{2} + c_2 = 0 \end{cases} \Rightarrow \begin{cases} c_2 = \frac{1}{2} \\ c_1 = \frac{1}{\sqrt{\pi}} \end{cases}$$

$$Q(x,t) = \frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{\frac{x}{\sqrt{4kt}}} e^{-\beta^2} d\beta$$

Now define $S = \frac{\partial Q}{\partial x}$ source function property b)

Remark: a) S solves the diffusion equation } We will see this later
 b) $\lim_{t \rightarrow 0^+} \frac{\partial Q}{\partial x}(t) = \delta$

First we observe that

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$$\partial_x Q(x,t) = \partial_x \left(\frac{1}{2} + \frac{1}{\sqrt{\pi}} \int_0^{\frac{x}{\sqrt{4tk}}} e^{-p^2} dp \right)$$

$$= \frac{1}{\sqrt{\pi}} e^{-\left(\frac{x}{\sqrt{4tk}}\right)^2} \cdot \frac{1}{\sqrt{4tk}} = S(x,t) \text{ source function!}$$

take a test function $\phi(x)$

$$\lim_{t \rightarrow 0} \left(\frac{1}{\sqrt{\pi}} e^{-\frac{x^2}{4t}} \frac{1}{\sqrt{4t}} \right) =$$

$$= \lim_{t \rightarrow 0} \text{so if } x \neq 0$$

$$\lim_{t \rightarrow 0^+} \frac{1}{\sqrt{\pi}} e^{-\frac{x^2}{4t}} \frac{1}{\sqrt{4t}} = 0$$

known exponential decays more than any polynomial

$$x = 0$$

$$\partial_x Q(0,t) = \frac{1}{\sqrt{\pi}} \cdot \frac{1}{\sqrt{4t}} \quad t \neq 0$$

so it looks singular, but if one uses the theory of distribution...

$$\left\{ \text{st. lim}_{|x| \rightarrow \infty} \phi(x) = 0 \right\}$$

Now assume that $\phi(x)$ is a given function & we define

$$u(x,t) = \int_{-\infty}^{\infty} S(x-y,t) \phi(y) dy \quad t > 0.$$

By (d) u is also a solution

Claim: u is the unique solution of

$$\begin{cases} u_t = \kappa u_{xx} \\ u|_{t=0}(x) = \phi(x) \quad \lim_{|x| \rightarrow \infty} \phi(x) = 0 \end{cases} \quad (4)$$

Proof: We only need to check the value of u at $t=0$

$$u(x,t) = \int_{-\infty}^{\infty} \mathcal{Q}(x-y, t) \phi(y) dy$$

$$= - \int_{-\infty}^{\infty} \mathcal{Q}_y(x-y, t) \phi(y) dy$$

$$= \underbrace{\mathcal{Q}(x-y, t) \phi(y)}_{\ddot{0}} \Big|_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \mathcal{Q}(x-y, t) \phi'(y) dy$$

So

$$u(x,0) = \int_{-\infty}^{\infty} \mathcal{Q}(x-y, 0) \phi'(y) dy$$

$$\parallel \begin{cases} 1 & \text{if } x-y > 0 \\ 0 & \text{if } x-y < 0 \end{cases}$$

$$= \int_{\{x-y > 0\}} \phi'(y) dy = \int_{-\infty}^x \phi'(y) dy = \phi(x)$$

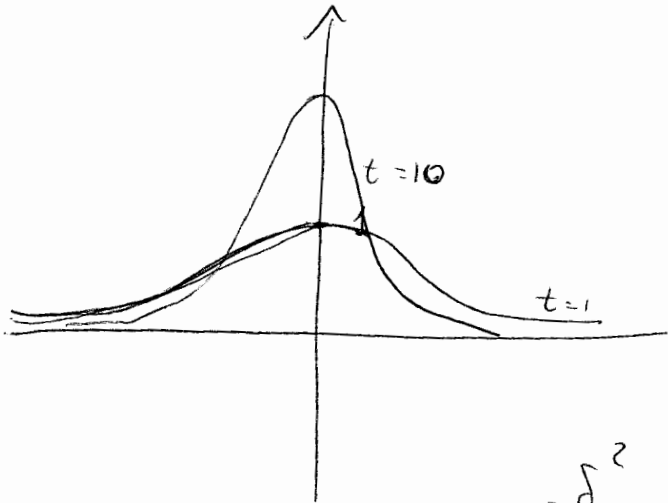
So

$$u(x,y) = \frac{1}{\sqrt{4\kappa\pi t}} \int_{-\infty}^{\infty} e^{-\frac{(x-y)^2}{4\kappa t}} \phi(y) dy$$

Q.E.D.

The function $S(x,t) = \frac{1}{2\sqrt{\pi kt}} e^{-x^2/4kt}$ $t > 0$

is called source function, Green's function, fundamental solution or propagator



$$f(x,t) = \frac{1}{\sqrt{t}} e^{-\frac{x^2}{t}}$$

$$f'(x,t) = \frac{1}{\sqrt{t}} e^{-\frac{x^2}{t}} \left(-\frac{2x}{t}\right)$$

$$t=1 \quad f(x,1) = e^{-x^2}$$

$$t=10$$

$$\max_{|x| > 5} S(x,t) = \frac{1}{\sqrt{t}} e^{-\frac{\delta^2}{t}}$$

$$\rightarrow t \rightarrow 0^+ \quad 0$$

Also $\int S(x,t) dx = \frac{1}{2\sqrt{\pi kt}} \int e^{-\frac{x^2}{4kt}} dx = 1$

This also supports the fact that $S(x,t) \rightarrow \delta_0$ as $t \rightarrow 0^+$

Problem: Solve $\begin{cases} u_t = k u_{xx} & k > 0 \\ u(x,0) = x^2 \end{cases}$

$$u(x,t) = \int \frac{1}{2\sqrt{\pi kt}} e^{-\frac{(x-y)^2}{4kt}} y^2 dy$$

$$\begin{aligned} \frac{x-y}{\sqrt{4kt}} &= p & dp &= -\frac{1}{\sqrt{4kt}} dy \\ & & &= \frac{1}{\sqrt{\pi}} \int dp e^{-p^2} (x - p\sqrt{4kt}) dp \end{aligned}$$

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$$= x^2 \int_{-\infty}^{\infty} \frac{1}{\sqrt{\pi}} e^{-p^2} dp + \frac{2x\sqrt{4kt}}{\sqrt{\pi}} \int e^{-p^2} p dp \quad (\star)$$

$$+ \frac{4kt}{\sqrt{\pi}} \int e^{-p^2} p^2 dp$$

$p^2 = \delta$ or $\delta = 2p dp \rightarrow$ this is not a trivial integral!

the way: u_{xxx} solves $u_t = k u_{xx}$ if u does (linearity)

$$u_{xxx}|_{t=0} = (x^2)_{xxx} = 0$$

so by uniqueness $u_{xxx}(x,t) = 0 \Rightarrow u_{xx}(x,t) = 2A(t)$

$$u_x(x,t) = 2A(t)x + B(t), \quad u(x,t) = A(t)x^2 + B(t)x + C(t)$$

$$\boxed{A'(t)x^2 + B'(t)x + C'(t) = k 2A(t)} \quad \text{equation } u_t = k u_{xx}$$

$$A'(t) = 0$$

$$B'(t) = 0$$

$$C'(t) = k 2A(t)$$

$$A(t) = c_0$$

$$B(t) = c_1$$

$$2k c_0 t + c_2 = C(t)$$

on the other hand

$$u(x,0) = A(0)x^2 + B(0)x + C(0) = x^2$$

$$B(0) = 0 \Rightarrow c_1 = 0 = B(t)$$

$$C(0) = 0 \Rightarrow c_2 = 0$$

$$A(0) = 1 \Rightarrow c_0 = 1$$

$$C(t) = 2kt \quad A(t) = 1$$

$$u(x,t) = x^2 + 2kt \quad (\star\star)$$

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For wave (*) with (**) we have

$$\frac{2}{\sqrt{\pi}} \int e^{-p^2} p^2 dp = 1$$

Comparison of wave and diffusion equation

Fundamental properties

Property:

	wave	Diffusion
i) Speed of propagation	Finite ($\leq c$)	Infinite
ii) Singularities for $t > 0$ (see initial data with 0 velocity)	Transported along characteristics (speed = c)	Lost immediately
iii) Well posedness for $t > 0$?	Yes	Yes (initial data)
iv) Well posedness for $t < 0$?	Yes	No
v) Maximum principle	No	Yes
vi) behaviour as $t \rightarrow \infty$	$E = \cos t$	$E \searrow$

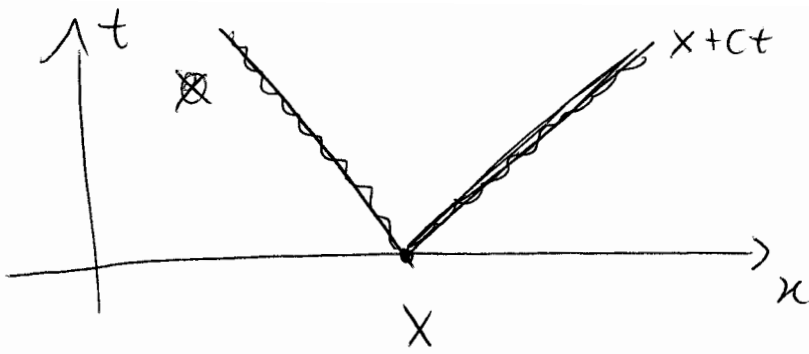
Some explanations

for wave:

if $\phi(x) = u(0, x)$, u solves $u_{tt} - c^2 u_{xx} = 0$

then we know that $u(x, t) = \frac{1}{2} [\phi(x+ct) + \phi(x-ct)]$

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for diffusion equation $u_t = k u_{xx}$ $u(x,0) = \phi(x)$

$$u(x,t) = \frac{1}{\sqrt{4\pi kt}} \int e^{-\frac{(x-y)^2}{4kt}} \phi(y) dy$$

all derivatives for u continuous on \mathbb{R}^2
 so if ϕ is in C^{∞} then u is C^{∞} but u is not!
 so ϕ maybe singular

so the value of u at (x,t) depends on the value of ϕ everywhere.

Similarly anything changed for ϕ at a certain point x_0 effects the values of u everywhere and at all times.

these facts take care of i) and ii)

iii) has been already discussed.

iv) if one changes variable $u(x,t) \rightarrow v(x,\tau)$

then $u(x,-\tau) = v(x,\tau)$

$v_{\tau\tau} = u_{tt}(x,-\tau)$ so also $v_{\tau\tau} - c^2 v_{xx} = 0$ $\tau > 0$
 $v_{xx} = u_{xx}(x,-\tau)$

This can be solved hence $u(x, t)$, $t < \infty$ can be solved.

On the other hand

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$$v_t = -u_{tt}(x, -c) \quad \text{so} \quad v_t = -k v_{xx} = 0$$

$$v_{xx} = u_{xx}(x, -c)$$

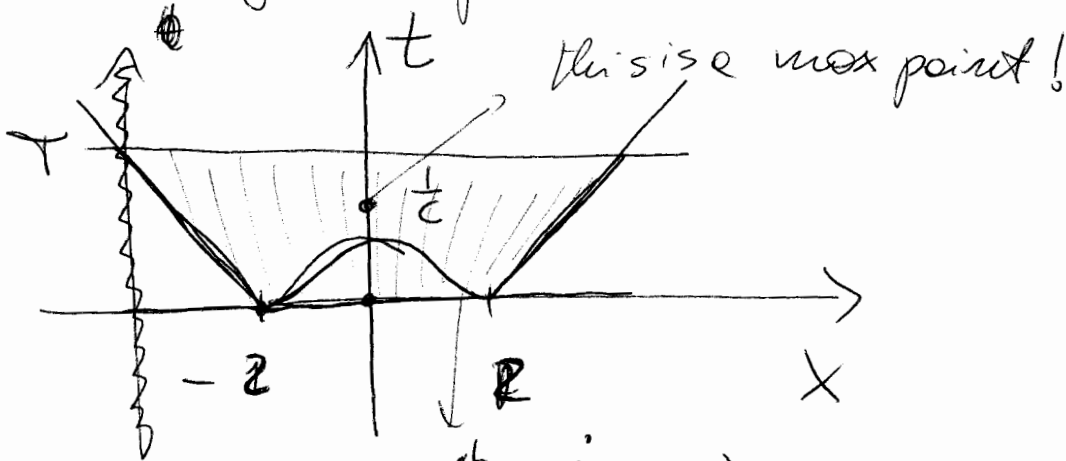
and one can see that the source function is now

$$\frac{1}{2\sqrt{\pi kt}} e^{-\frac{x^2}{4kt}}$$

which is not integrable $c > 0$

not integrable.

v) that this is not true for the wave comes from the following example



$$\phi = u(0, x)$$

$$\phi = 0$$

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} \psi(s) ds$$

$$x_0 = 0$$

$$ct = 2 \quad t = \frac{1}{c}$$

Summarizing:

- 1) Classification of ODE $\left\{ \begin{array}{l} \text{linearity, homogeneous, inhomogeneity} \\ \text{non linearity} \\ \text{order} \end{array} \right.$
- 2) Solution to ^{linear} 1st order PDE in \mathbb{R}^2 (homogeneous + inhomogeneous)
- 3) Classification of 2nd order linear PDE $\left\{ \begin{array}{l} \text{elliptic} \\ \text{Hyperbolic} \\ \text{parabolic} \end{array} \right.$
 - Types
 - 2 initial and boundary conditions
- 3) Well-posedness,
- 6) Wave equation in \mathbb{R}^1 $\left\{ \begin{array}{l} \text{solution of initial value problem} \\ \text{causality} \\ \text{Energy} \\ \text{Source functions} \end{array} \right.$
- 7) Diffusion Equation $\left\{ \begin{array}{l} \text{Maximum principle} \\ \text{uniqueness} \\ \text{Energy} \\ \text{Source functions} \\ \text{Solution of the IVP} \\ \text{on } \mathbb{R}. \end{array} \right.$